

Light Quanta on Beam Splitters

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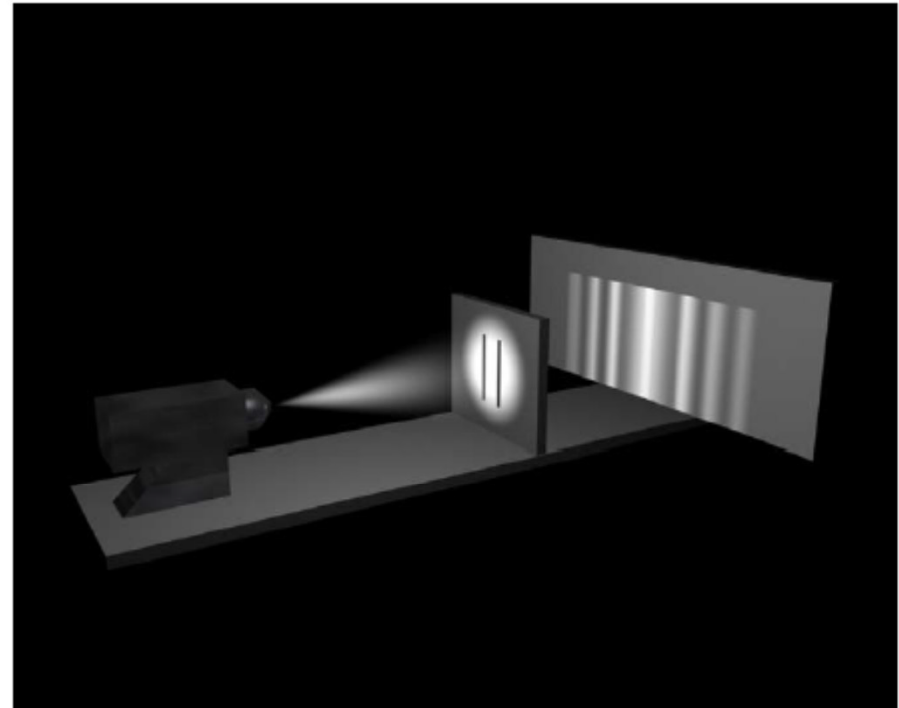
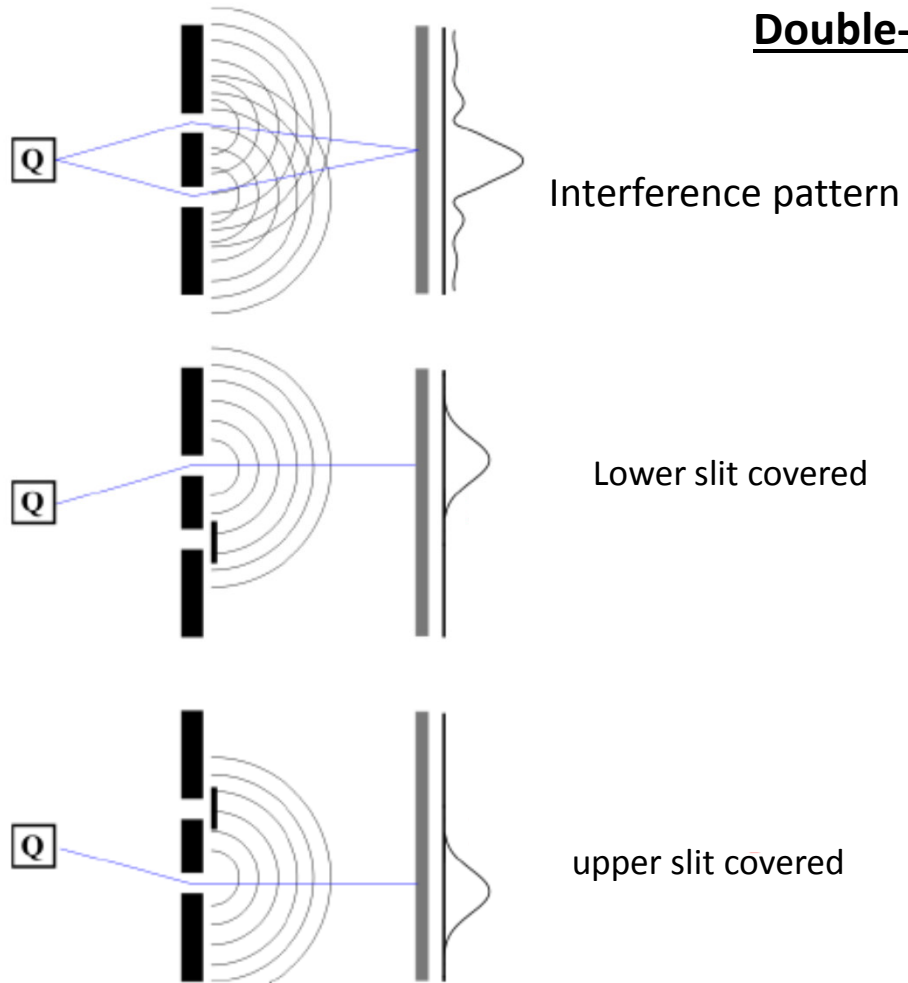
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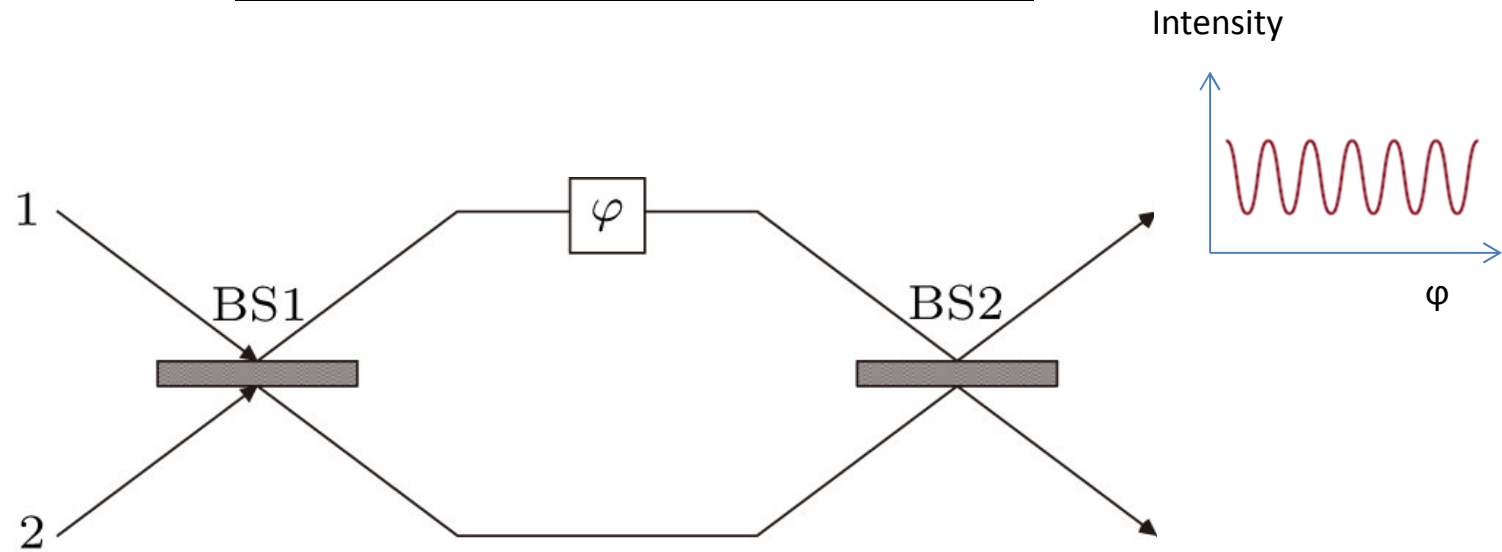
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1. Introduction

Double-slit experiment



Mach-Zehnder Interferometer



BS1 = 1. Beam Splitter

BS2 = 2. Beam Splitter

φ = Phase Shifter

2. Quantum harmonic oscillator, photons and Wigner function

QM harmonic oscillator

$$\hat{H} = \frac{1}{2} \left[\frac{1}{m} \hat{p}^2(t) + m\omega^2 \hat{q}^2(t) \right] = \hbar\omega \left(\hat{N} + \frac{1}{2} \hat{1} \right), \quad \hat{N} = \hat{a}^+ \hat{a}, \quad \hat{N}|n\rangle = n|n\rangle$$

$$\hat{H}|n\rangle = E_n|n\rangle \quad \rightarrow$$

$$\langle q|n\rangle = \psi_n(q) \propto H_n \left(\frac{q}{q_0} \right) \exp \left(-\frac{q^2}{2q_0^2} \right)$$

$$\left. \begin{aligned} \hat{a} &= \frac{1}{\sqrt{2\hbar}} \left(\lambda \hat{q} + \frac{i}{\lambda} \hat{p} \right) \\ \hat{a}^+ &= \frac{1}{\sqrt{2\hbar}} \left(\lambda \hat{q} - \frac{i}{\lambda} \hat{p} \right) \end{aligned} \right| \lambda = \sqrt{m\omega}$$

$$q_0 = \sqrt{\frac{\hbar}{m\omega}}$$

Hermite polynomials

$|n\rangle$ = photon state
or light quantum
($n=0,1,2,\dots$)

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right),$$

For light $m := \hbar\omega/c^2$

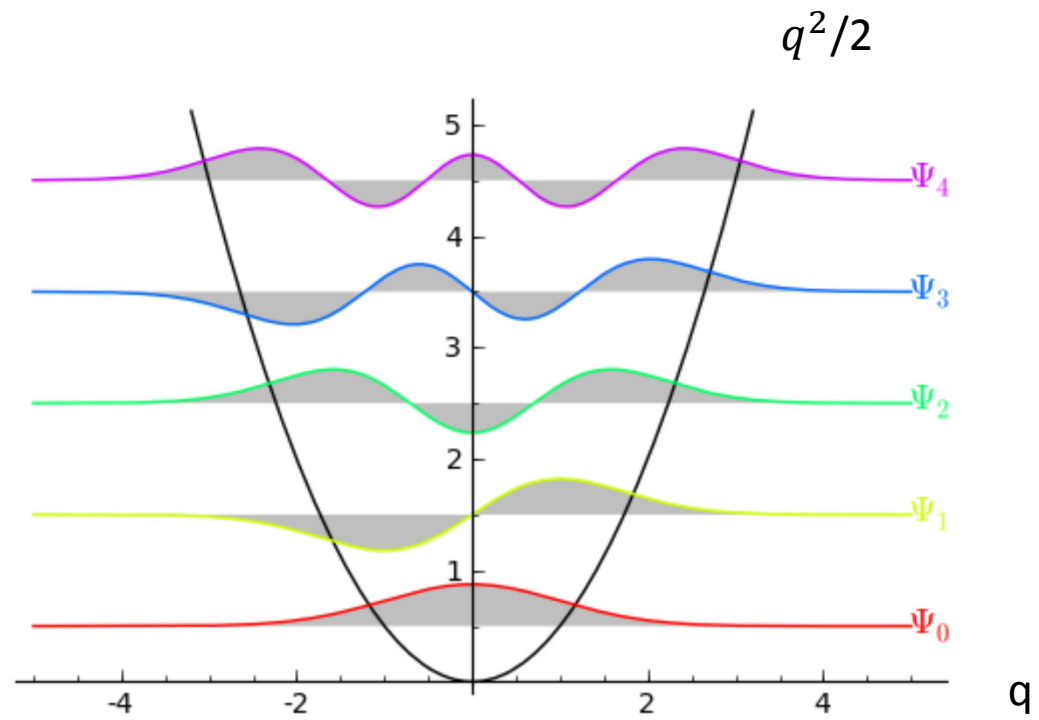
$W(q,p)$ = Wigner function

q = position

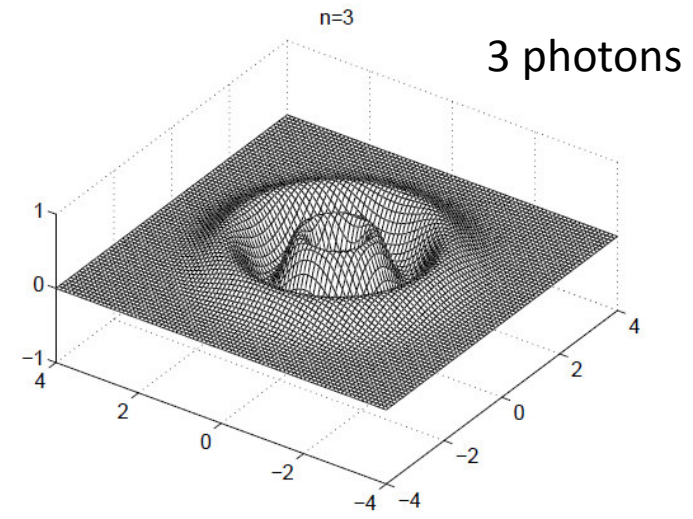
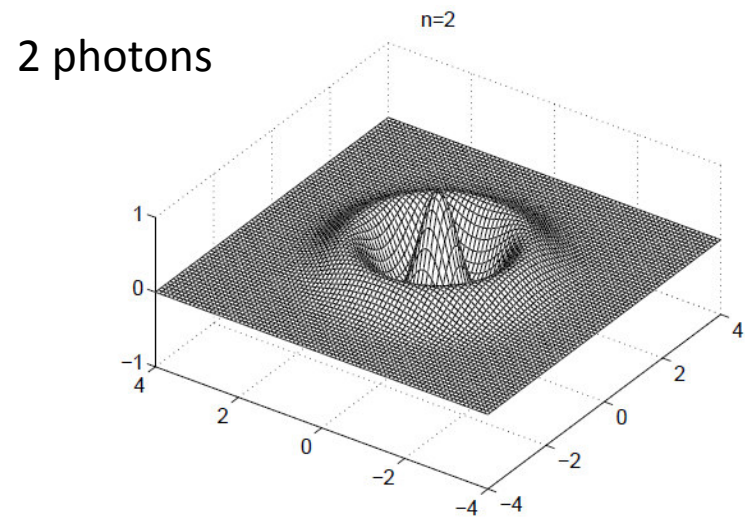
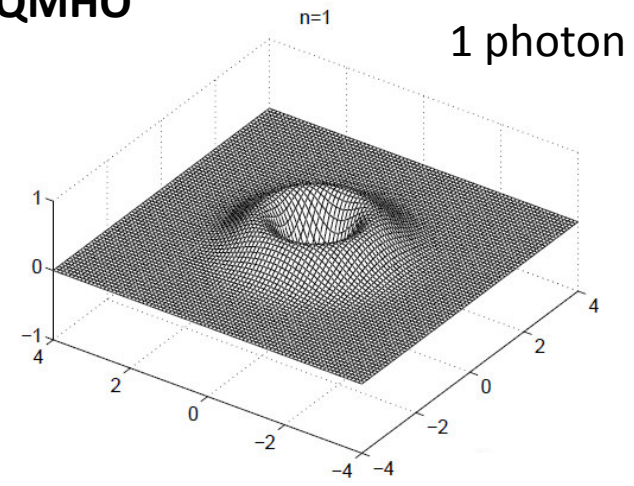
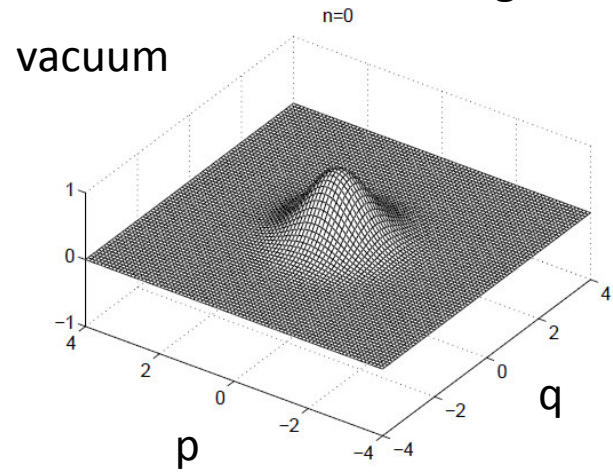
p = momentum

$$W(q,p) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} \psi_n \left(q - \frac{y}{2} \right) \psi_n^* \left(q + \frac{y}{2} \right) e^{iyp/\hbar} dy$$

QMHO-Potential

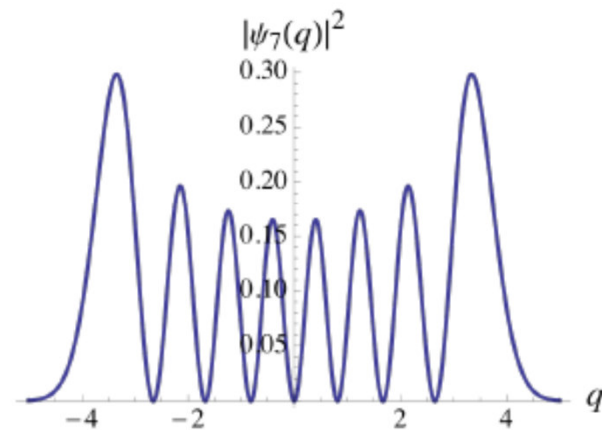


Wigner function of QMHO

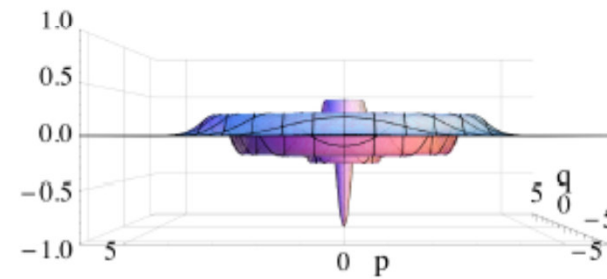
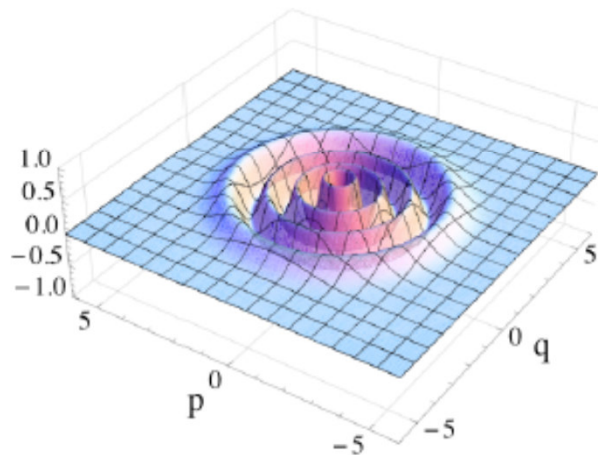


QM harmonic oscillator

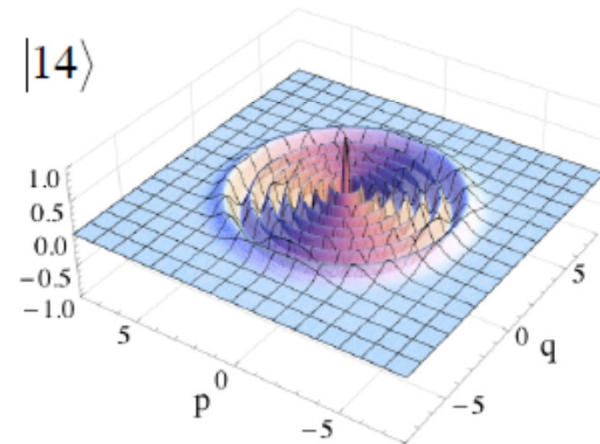
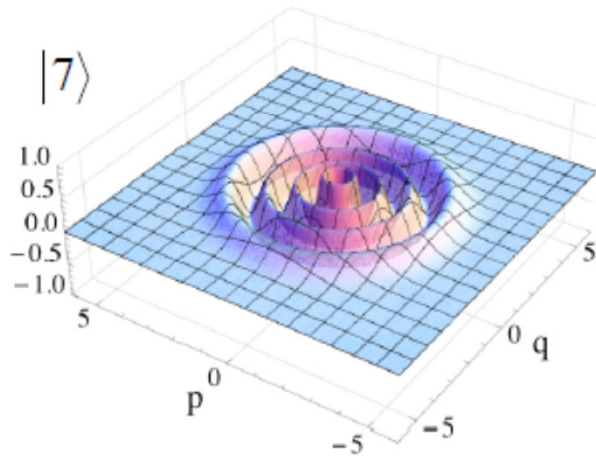
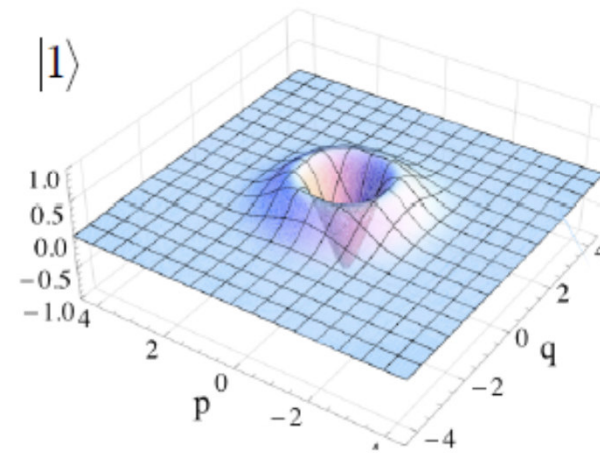
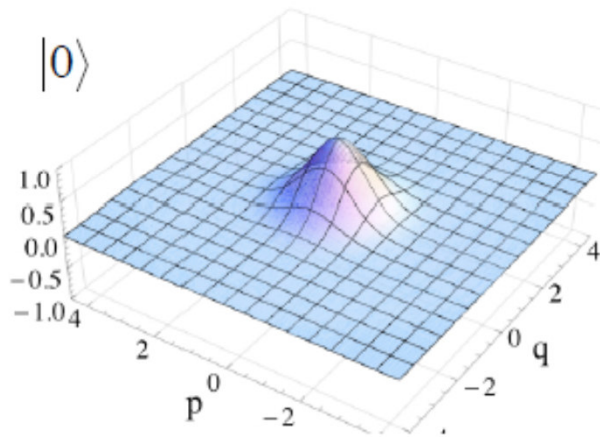
$|n = 7\rangle$



$\hbar = 1,$
 $q_0 = 1,$
 $W \times \pi.$

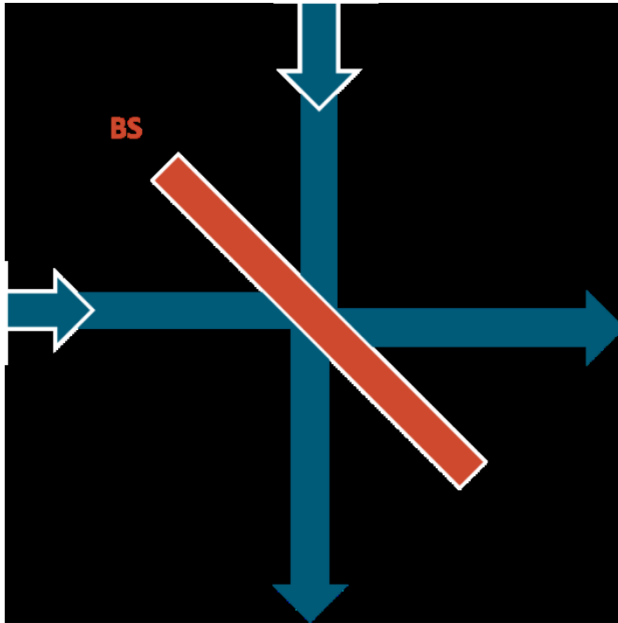


Fock states: Wigner function



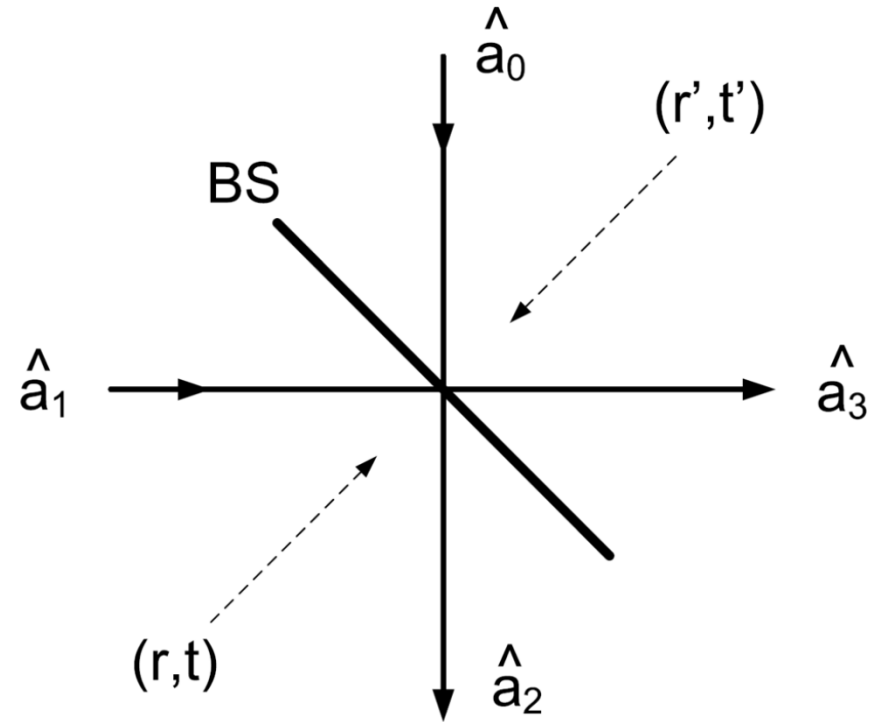
3. Beam splitter

Glass – Reflection and Transmission



We choose for reflection coefficient: $r = i/\sqrt{2}$, $i = \exp(i\pi/2)$
 Transmission coefficient: $t = 1/\sqrt{2}$

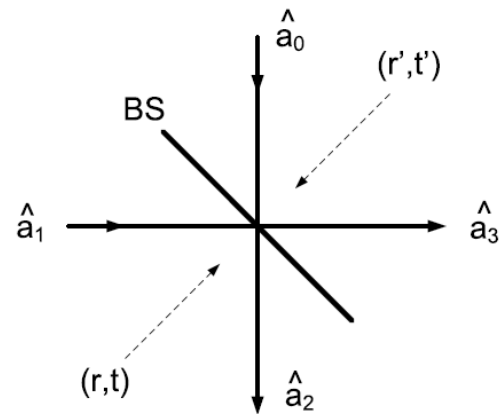
Simplified & Balanced Beam Splitter



$$\hat{a}_2 = t\hat{a}_0 + r\hat{a}_1$$

$$\hat{a}_3 = r\hat{a}_0 + t\hat{a}_1$$

50:50 Beam splitter BS or \hat{T}



a)
$$\begin{pmatrix} \hat{a}_2 \\ \hat{a}_3 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & i \\ i & 1 \end{pmatrix} \begin{pmatrix} \hat{a}_0 \\ \hat{a}_1 \end{pmatrix} = \hat{T} \begin{pmatrix} \hat{a}_0 \\ \hat{a}_1 \end{pmatrix}$$
 with BS-operator \hat{T}

b) Alternative description of the BS:

$$\begin{pmatrix} \hat{a}_2 \\ \hat{a}_3 \end{pmatrix} = \hat{U}^\dagger \begin{pmatrix} \hat{a}_0 \\ \hat{a}_1 \end{pmatrix} \hat{U}$$

with unitary operator

$$\hat{U} = e^{-i\hat{H}\tau}$$

Hamiltonian of BS:

$$\hat{H} = \kappa (\hat{a}_0^\dagger \hat{a}_1 + \hat{a}_0 \hat{a}_1^\dagger) \quad , \kappa = 1/2, \tau = \pi/2$$

\hat{a}^\dagger and \hat{a} are creation and annihilation operators

$$\hat{a}|n\rangle = \sqrt{n}|n-1\rangle, \hat{a}^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle \quad \longrightarrow \quad \begin{aligned} \hat{a}^\dagger|0\rangle &= |1\rangle \\ \hat{a}|1\rangle &= |0\rangle \end{aligned}$$

4. Coherent state

Coherent state (CS) $|\alpha\rangle$

BS not only for Fock-states $|n\rangle$, but also with CS $|\alpha\rangle$:

$$|\alpha\rangle = \hat{D}(\alpha)|0\rangle = e^{-|\alpha|^2/2} e^{\alpha\hat{a}^+} |0\rangle =$$

$$= e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{\alpha^n}{n!} (\hat{a}^+)^n |0\rangle = e^{-|\alpha|^2/2} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle \quad , \quad \hat{D}(\alpha) = \text{displacement-operator}$$

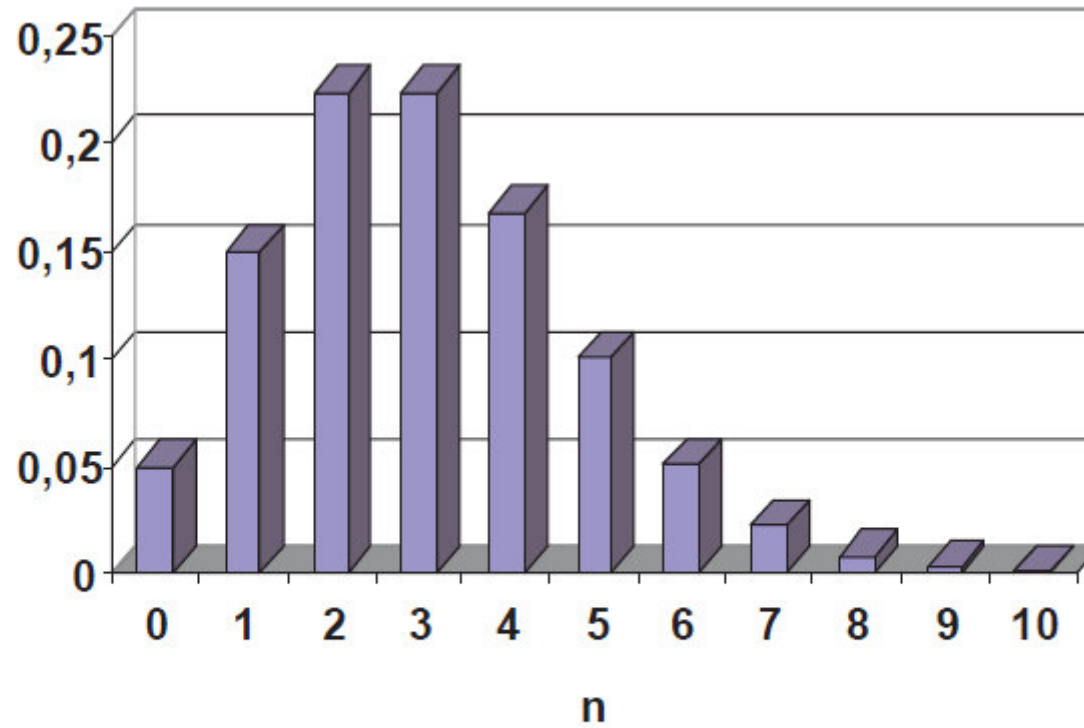
$$\approx |0\rangle + \alpha |1\rangle \quad , \quad |\alpha|^2 \ll 1$$

$$|\alpha|^2 = \text{mean photon number}$$

$$\hat{D}(\alpha) = e^{\alpha\hat{a}^+ - \alpha^*\hat{a}} \quad ,$$

$$\alpha \in \mathbb{C}$$

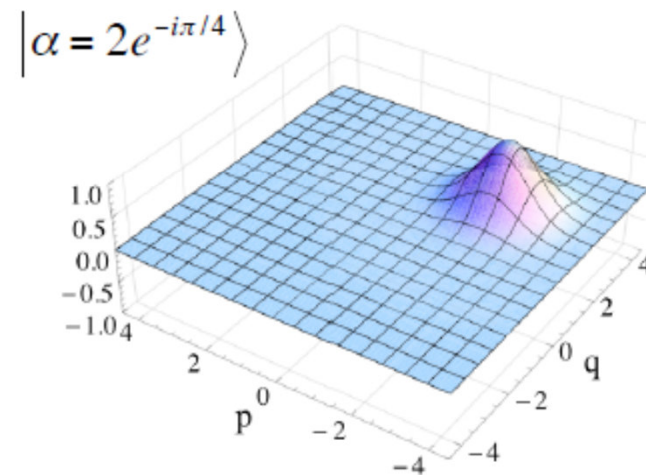
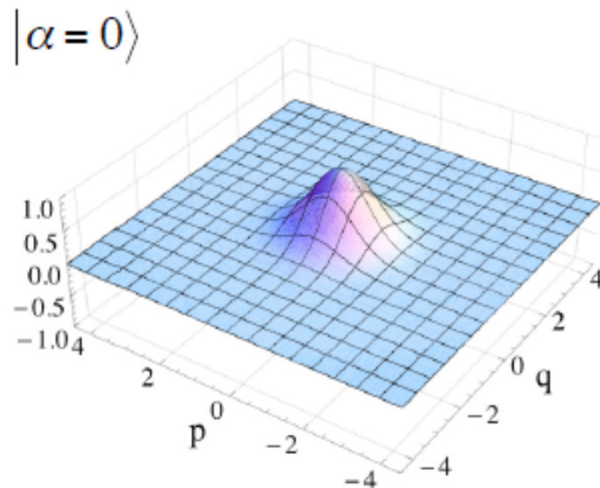
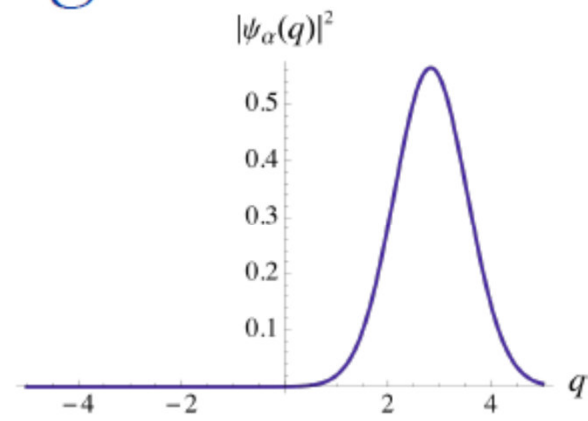
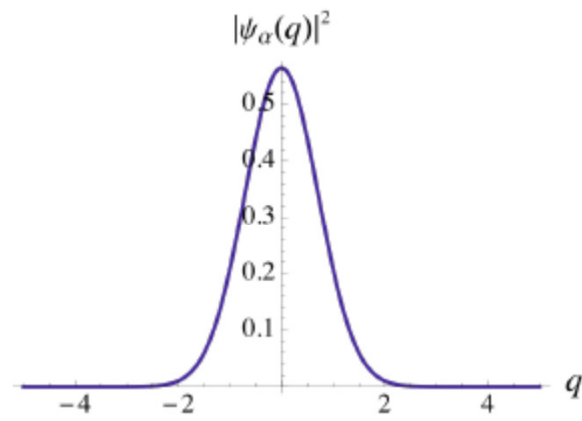
$$\langle\alpha|\alpha\rangle = e^{-|\alpha|^2} \sum_{n=0}^{\infty} \frac{|\alpha|^{2n}}{n!} = \sum_{n=0}^{\infty} P_n = 1$$



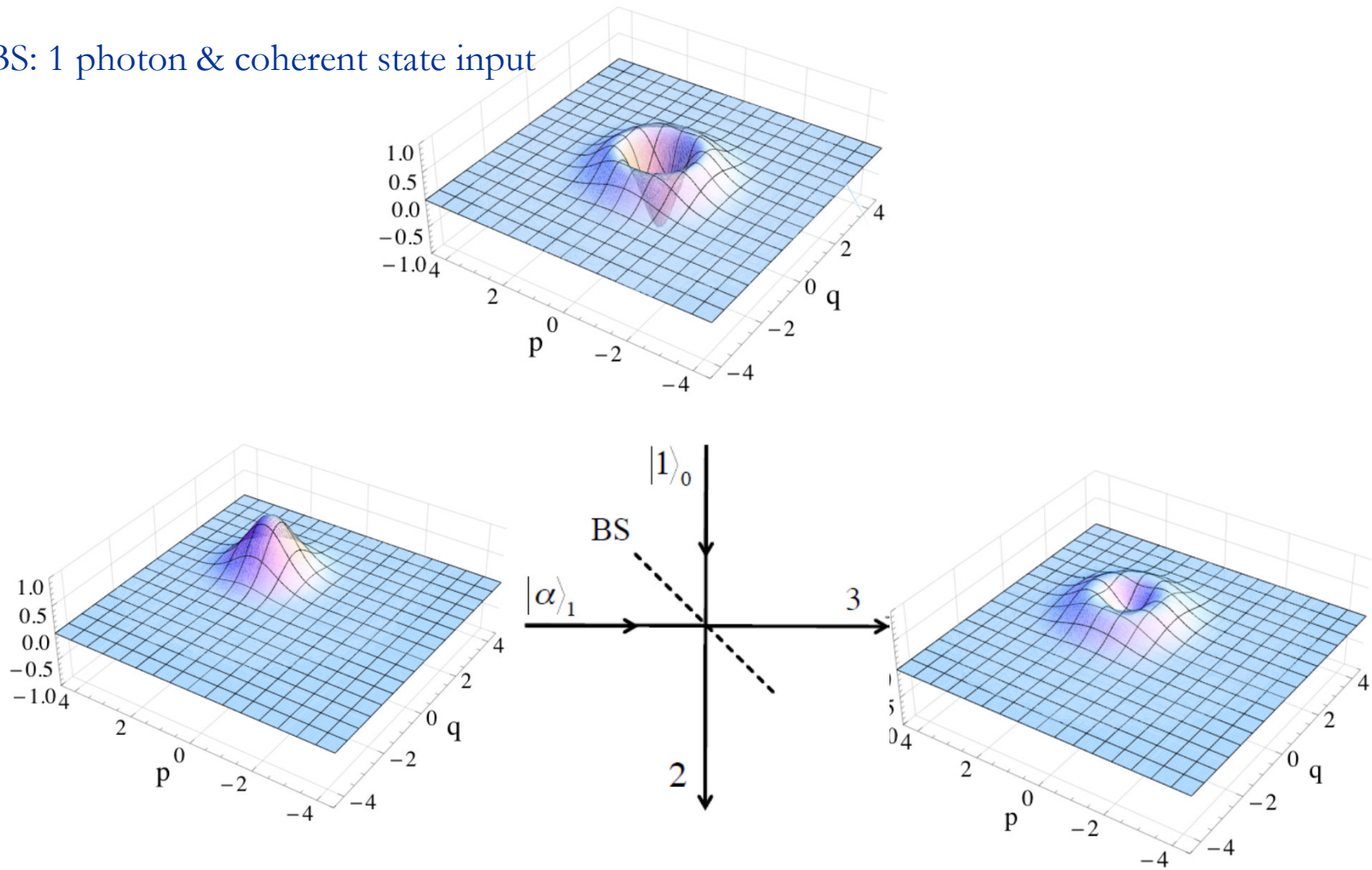
Poissonian distribution $P_n = e^{-|\alpha|^2} \frac{|\alpha|^{2n}}{n!}$ for $\langle n \rangle = |\alpha|^2 = 3$

$|\alpha|^2 = \text{mean photon number of CS}$

Coherent states: Wigner function



BS: 1 photon & coherent state input



5. Quantum random numbers (QRN): Examples

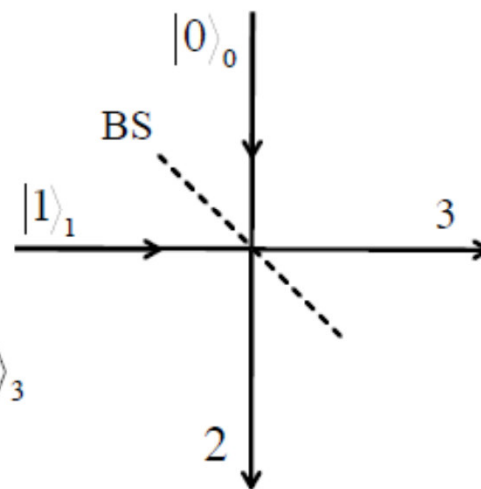
1. Example:

Beam splitter: 1 photon input

$$|0\rangle_0|0\rangle_1 \xrightarrow{BS} |0\rangle_2|0\rangle_3 \quad a^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle$$

$$|0\rangle_0|1\rangle_1 = \hat{a}_1^\dagger|0\rangle_0|0\rangle_1 \quad \hat{a}_1^\dagger = r\hat{a}_2^\dagger + t\hat{a}_3^\dagger$$

$$\hat{a}_1^\dagger|0\rangle_0|0\rangle_1 \xrightarrow{BS} (r\hat{a}_2^\dagger + t\hat{a}_3^\dagger)|0\rangle_2|0\rangle_3 = r|1\rangle_2|0\rangle_3 + t|0\rangle_2|1\rangle_3$$



Balanced Beam Splitter equation:

A single input photon in mode 1 together with a vacuum input in mode 0 is equally transmitted and/or reflected with a probability of $\frac{1}{2}$:

$$\begin{aligned} |0\rangle_0 |1\rangle_1 &\xrightarrow{BS} \frac{1}{\sqrt{2}} (i\hat{a}_2^\dagger + \hat{a}_3^\dagger) |0\rangle_2 |0\rangle_3 = \\ &= \frac{1}{\sqrt{2}} (i|1\rangle_2 |0\rangle_3 + |0\rangle_2 |1\rangle_3) . \end{aligned} \quad \Rightarrow \text{QRN !}$$

Remember:

We choose for reflection coefficient: $r=i/\sqrt{2}$, $i=\exp(i\pi/2)$

Transmission coefficient : $t = 1/\sqrt{2}$

2. Example:

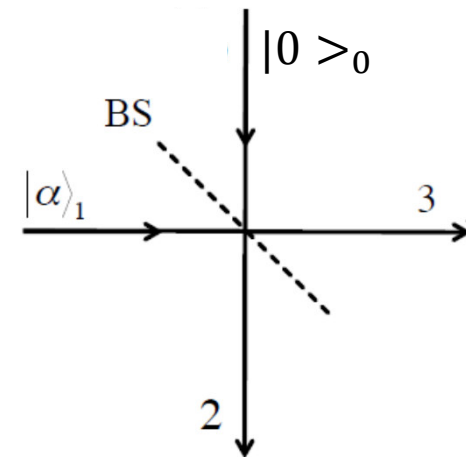
A CS contains an infinite number of Fock states $|n\rangle$.
 For small mean photon number $|\alpha|^2 \approx 1/100$ we
 obtain quantum random numbers (*cf. paper*)

Input: $|0\rangle_0 |\alpha\rangle_1$

Output:

$$\left| \frac{i\alpha}{\sqrt{2}} \right\rangle_2 \left| \frac{\alpha}{\sqrt{2}} \right\rangle_3 \approx |0\rangle_2 |0\rangle_3 + \frac{\alpha}{\sqrt{2}} [i|1\rangle_2 |0\rangle_3 + |0\rangle_2 |1\rangle_3]$$

Here $|\alpha\rangle \approx |0\rangle + \alpha|1\rangle$, $|\alpha|^2 \ll 1$ has been used.



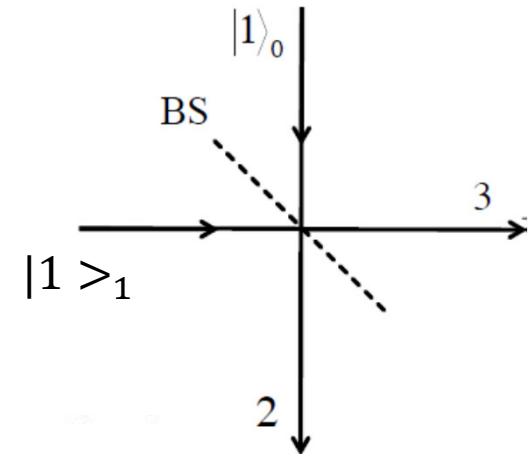
➔ QRN !

3. Example:

Hong Ou Mandel (*HOM*) effect; Input of 2 single photons

Result: 2 photons with probability $\frac{1}{2}$ on either output-mode 2 or 3

$$\begin{aligned} |1\rangle_0 |1\rangle_1 &= \hat{a}_0^+ \hat{a}_1^+ |0\rangle_0 |0\rangle_1 \xrightarrow{BS} \rightarrow \\ &= \frac{i}{\sqrt{2}} (|2\rangle_2 |0\rangle_3 + |0\rangle_2 |2\rangle_3) \end{aligned}$$

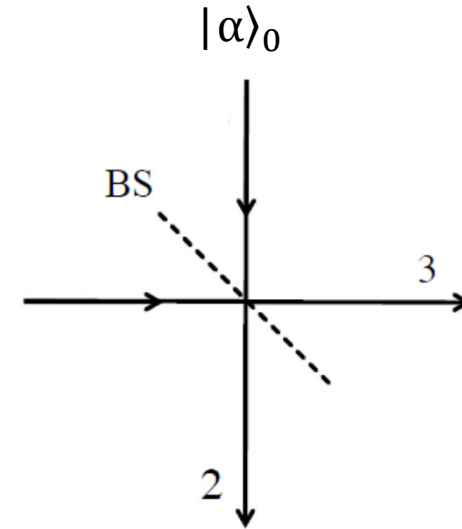


➡ QRN !

4. Example:

CS on input 0 and Schrödinger-cat-state on input 1

$$(|\alpha\rangle_1 + |-\alpha\rangle_1) / N$$



Output:

$$|\psi\rangle_{out} = \frac{1}{N} [|\sqrt{2}\alpha\rangle_2 |0\rangle_3 + |0\rangle_2 |\sqrt{2}\alpha\rangle_3] ,$$

$$N = \sqrt{2(1 + e^{-2|\alpha|^2})} .$$

➡ QRN !

6. Summary

- Double split experiment, Mach-Zehnder interferometer
- Quantum harmonic oscillator, photons and Wigner function
- Beam splitter
- Coherent state
- Quantum random numbers (QRN): 4 Examples

Thank you very much!